

GRAVITATIONAL WAVE GENERATION FROM COSMOLOGICAL PHASE TRANSITIONS

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3^o Edition of the Friends-of-Friends Meeting
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- 2 The electroweak phase transition
 - Bubble nucleation
 - Hydrodynamics of the phase transition
- 3 Turbulent source of gravitational waves
 - Hydrodynamic parameters
 - Gravitational waves
- 4 Work in progress

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What are we interested in?

Particle Physics predicts several phase transitions:

- Grand Unification (GUT) ($T \sim 10^{15}\text{GeV}$)
- **Electroweak Transition** ($T \sim 100\text{GeV}$)
- Quark-Hadron Transition (QCD) ($T \sim 100\text{MeV}$)

(We use *natural units*: $c = \hbar = k_B = 1$ and $\text{eV} \sim 10^4\text{K}$)

This talk is focused on the electroweak phase transition, which has several potentially observable cosmological consequences:

- Magnetic fields
- Baryogenesis, baryonic inhomogeneities
- **Gravitational Waves** (This is what interests us!)

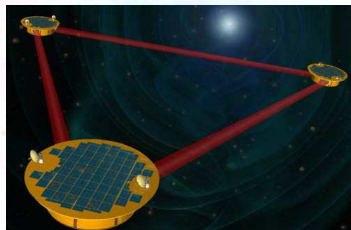
Note:

Most of the concepts that will be discussed here are valid for other phase transitions.

Gravitational Wave generation *in a phase transition*

Bulk motions of the relativistic fluid in a phase transition are a source of gravitational waves.

- Unlike electromagnetic radiation, gravitational radiation, once generated, propagates freely.
 - It could provide information about earlier stages of the universe.
- The frequency of gravitational waves generated in the electroweak phase transition would be \sim **miliHertz**.
 - It lies within **LISA**'s sensibility range.



Laser Interferometer
Space Antenna

Characteristic frequency and amplitude

The spectrum essentially depends on a few parameters related with the dynamics of the transition

- The peak wavelength of the gravitational radiation is determined by the **characteristic length** of the source , $\lambda_p \sim L_S$
- Then, peak frequency is given by $f_p \sim 1/L_S$
- Once generation occurs, the waves propagate freely until present
- Taking into account the redshift, the peak frequency today is given by

$$f_p \sim \frac{1}{L_S} \left(\frac{a_*}{a_0} \right) \sim 10^{-5} \text{ Hz} \frac{T}{100 \text{ GeV}} \frac{H_*^{-1}}{L_S}.$$

- a_* and a_0 are the scale factor at that moment and today.
- 100 GeV = electroweak scale
- H_*^{-1} = Hubble length at that moment ($H = \dot{a}/a$ Hubble rate)
- (Here, we used adiabatic expansion and a Hubble rate given by the Friedmann equation.)

- Considering a perturbation $h_{\mu\nu}$ of the *metric* $g_{\mu\nu}$, linearizing *Einstein equation* and keeping the relevant degrees of freedom, we can obtain a *wave equation* of the form: $\square h \sim GT$

- \square is a linear operator which contains second derivatives.
- T = A projection of the energy-momentum tensor of the source.
- G = Newton constant.

- The length scale of the source gives the variation scale, $\partial_\mu \sim \frac{1}{L_S}$

We can estimate h dimensionally: $L_S^{-2} h \sim G \rho_K \Rightarrow h \sim G \rho_K L_S^2$

- ρ_K = kinetic energy density in bulk motions of fluid.

- The gravitational waves intensity is closely related to energy density:

$$\rho_{\text{GW}}(\mathbf{x}, t) \sim \langle \partial_t h_{\mu\nu} \partial_t h^{\mu\nu} \rangle / G \sim (\partial_t h)^2 / G \sim (L_S^{-1} h)^2 / G$$

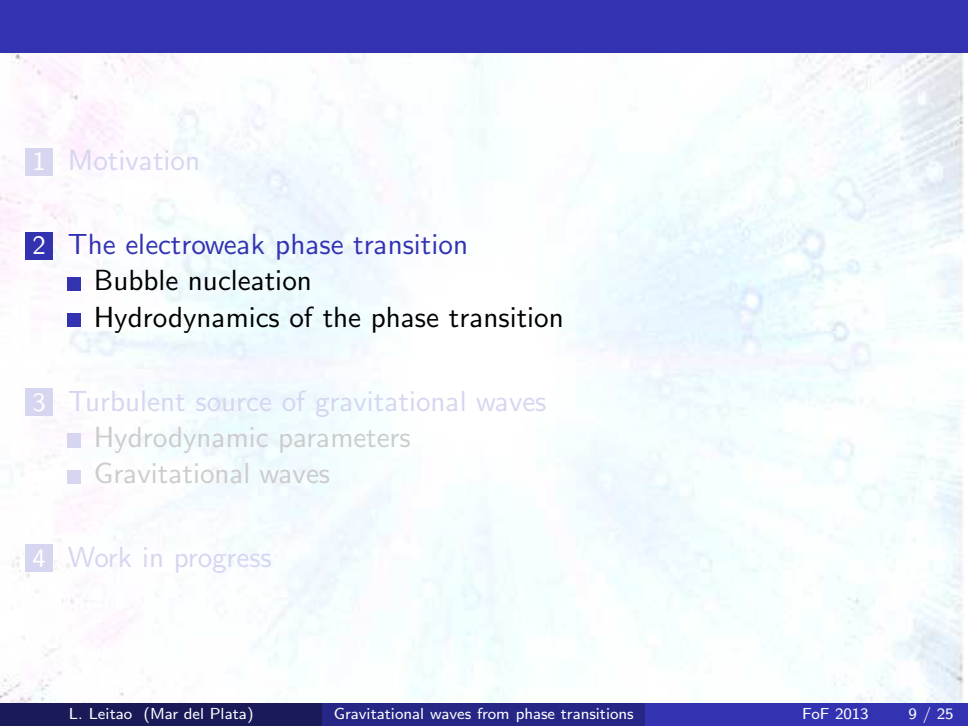
- Then, the wave energy density is $\rho_{\text{GW}} \sim G \rho_K^2 L_S^2$

- Taking into account the redshift, the energy density **today** is

$$\rho_{\text{GW}} \sim (\rho_K / \rho_*)^2 (L_S H_*)^2 \rho_0$$

- (using $H_*^2 \sim G \rho_*$, *Friedmann Equation*)
- ρ_* total energy density at this moment
- $\rho_0 = \rho_* (a_*/a_0)^4$ radiation energy density today

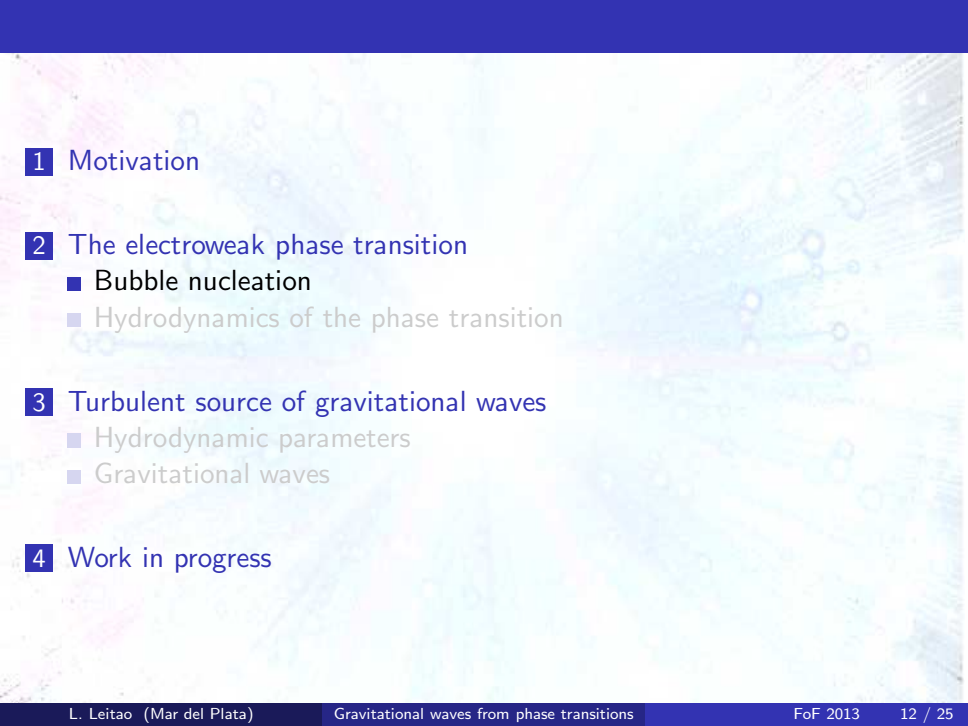
- We are interested in gravitational wave generation from the electroweak phase transition.
 - Expected signal:
 - $\rho_{\text{GW}} \sim \rho_{\text{K}}^2 L_S^2$
 - $f \sim \frac{T_*}{L_S}$
- (Peak value estimations)
- In order to estimate the relevant parameters, it is necessary to study the hydrodynamics involved.

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- Part of the particle physics *Standard Model*.
- Unifies electromagnetic and weak nuclear interactions.
- Involves all *Standard Model* particles: quarks, leptons, photons, Z & W bosons, and the Higgs field.
- The particle masses depend on the Higgs field.

What is “the electroweak phase transition”?

- **Reminding thermodynamics:**
 - Every phase transition is governed by a thermodynamic potential (free energy) which depends on an order parameter.
 - In different regions of the universe, the order parameter takes different values, corresponding to different free-energy minima.
- In the electroweak phase transition the order parameter is the background (thermodynamic average) of the **Higgs field**.
- During the transition there is a coexistence of two minima associated with two phases. Each phase is characterized by:
 - Radiation (massless particles) - dominates at high temperature.
 - Massive particles (due to Higgs field) - dominates at low temperature.

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Bubble nucleation

Classification of phase transitions:

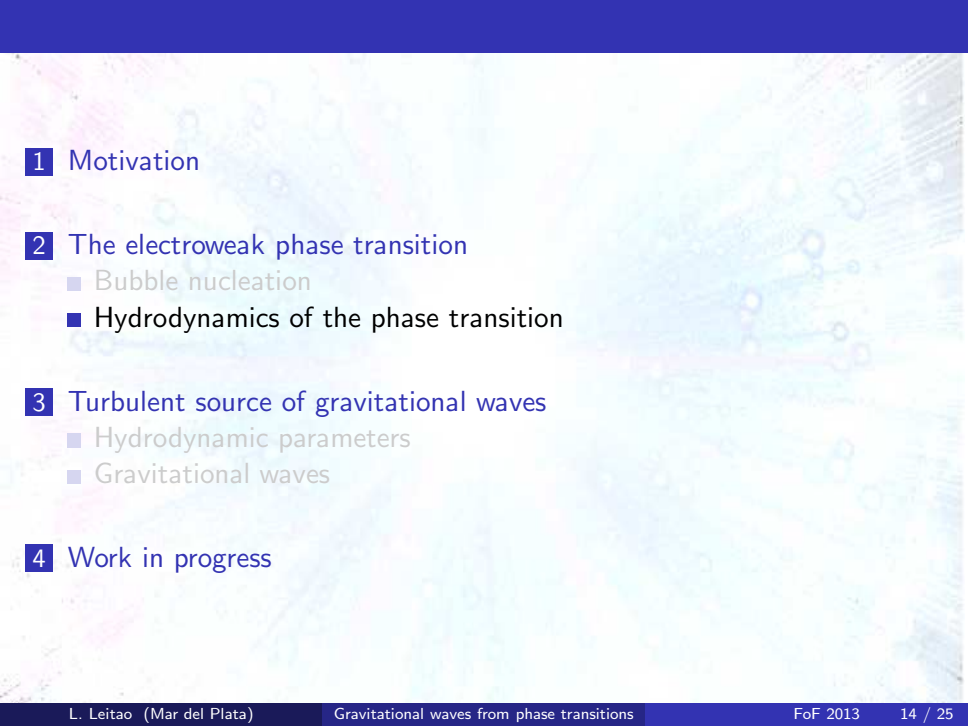
- **1st order:** Supercooling or superheating occurs, there is a latent heat (an energy that, when is released, can “shake” and reheat the medium).
- **Superior order:** There is not a latent heat (Very boring!).

Characteristics of a strongly 1st order phase transition:

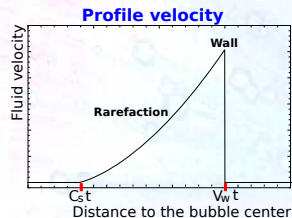
- Domains (bubbles) are nucleated containing the phase that is dominant at low temperature.
- Due to the supercooling there is a pressure difference between phases, which causes bubble expansion.
- The stronger the supercooling, the more violent the expansion and the stirring of the medium. **This causes high-intensity gravitational waves.**

Note 1: The 1st order transitions are not only interesting for us, also for those who study baryogenesis.

Note 2: Different extensions of the particle physics Standard Model may have different transition strength and hydrodynamics in the bubbles.

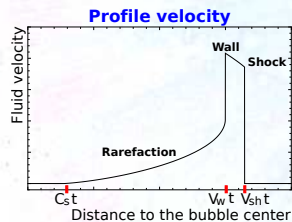
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The bubble wall has two different propagation modes:



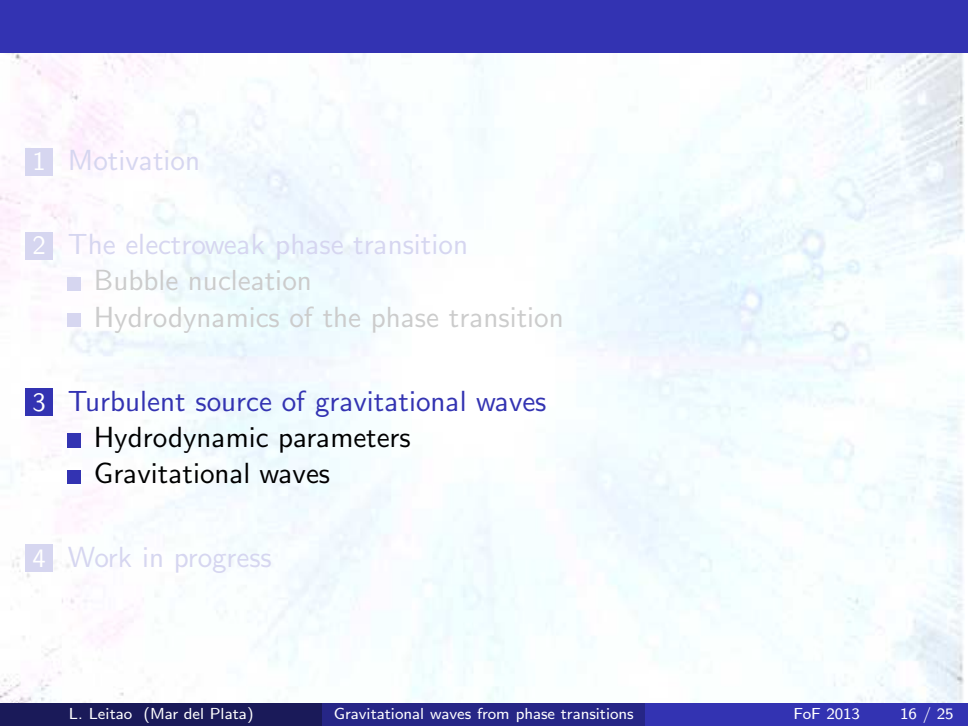
Detonation

- Supersonic wall (relativistic gas, $c_s \approx 1/\sqrt{3}$)
- Unperturbed fluid in front the wall
- *Rarefaction wave* following behind the wall



Deflagration

- Subsonic wall (respect to adjacent fluid)
- *Shock wave* preceding the wall
- Unperturbed fluid behind the wall (Usually)
Although if the wall is supersonic, there is also a rarefaction "tail"

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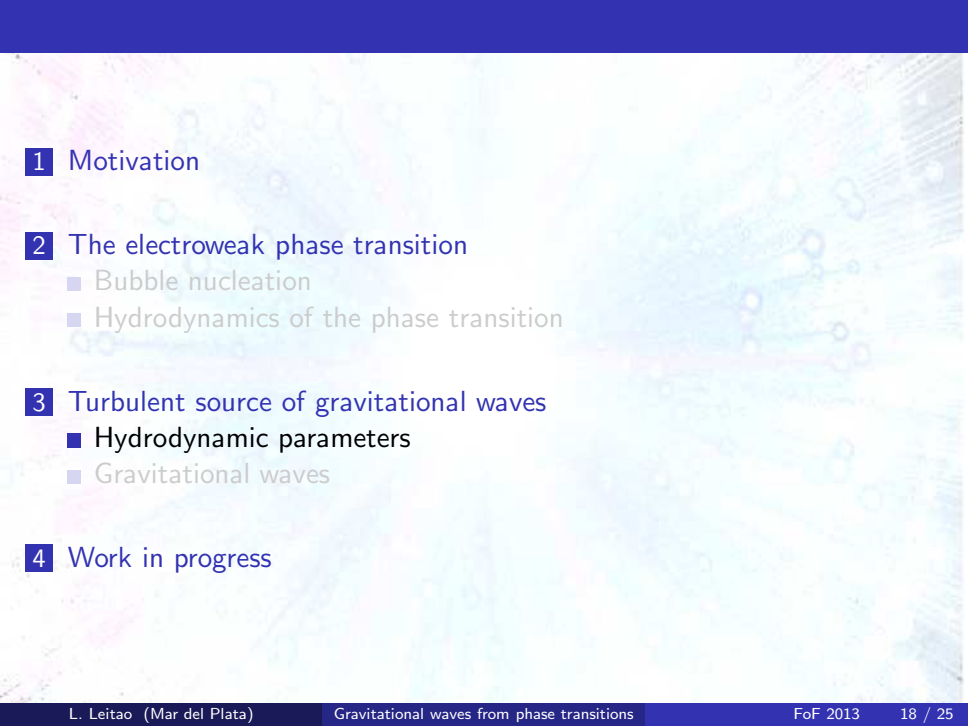
Turbulence in a phase transition

- The nucleated bubbles expand, either as detonations or deflagrations, dragging or pushing fluid until they collide generating turbulence.
- The generated eddies grow and also break into smaller eddies.
- So, there are different scales of eddies, which are source of gravitational waves with different wavelength.

Some calculation assumptions

This subject is complex and has several approximations. . .

- The phase transition is injecting energy at a single length scale, although there are bubbles with different sizes.
- There is some ambiguity to decide which scale is most relevant.
- In general, the literature in this subject doesn't consider the entire process from bubble collisions to the state of fully developed turbulence.
(Great!, we have job yet)

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Determination of hydrodynamic parameters

Remember these parametric dependencies:

- $\rho_{\text{GW}} \sim \rho_K^2 L_S^2$
- $f \sim \frac{T_*}{L_S}$

By studying the thermodynamics of the phase transition we can calculate:

- The temperature at which bubble nucleation begins T_* .
- The latent heat.

In conjunction with the hydrodynamics study, we know at any time:

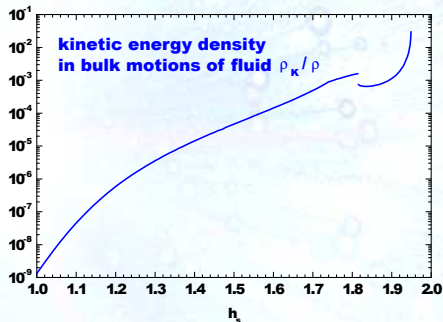
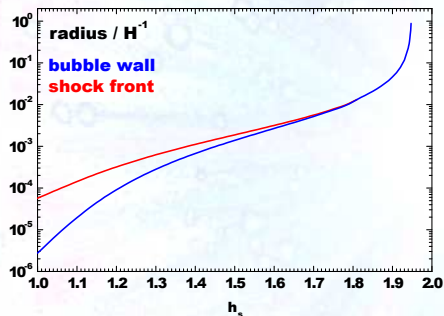
- The bubble diameter as a function of its nucleation time (or temperature).
- The fluid velocity profile around the wall of each bubble.
(Hence, the profile of kinetic energy distribution)

Assuming that turbulence *inherits* features present at collision time, the size of the biggest bubbles and the average kinetic energy density can be used to estimate the generation parameters L_S and ρ_K .

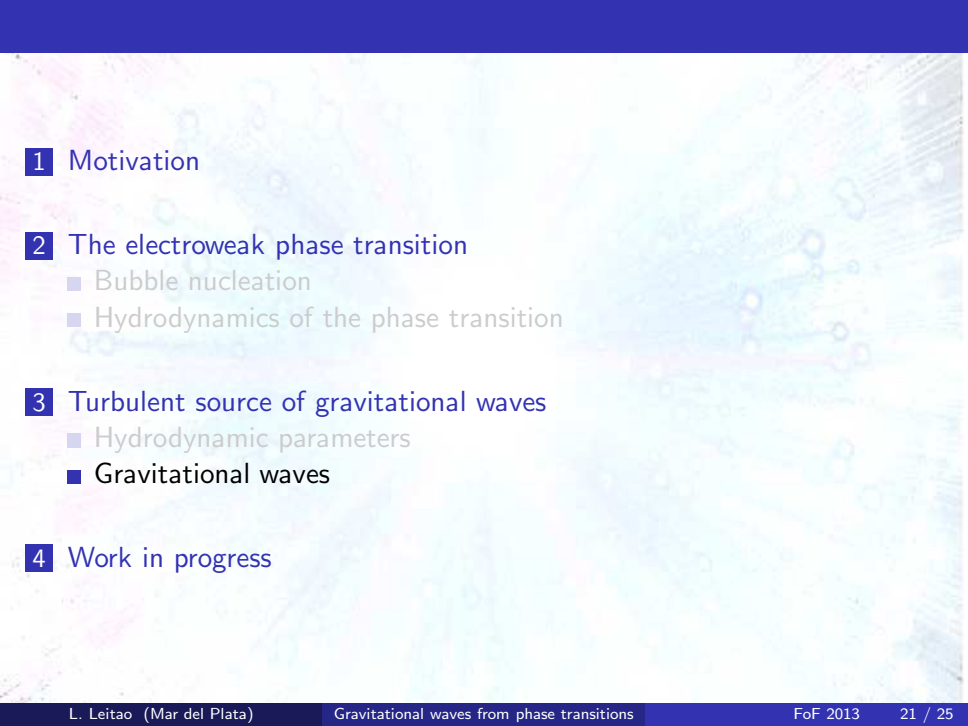
Hydrodynamic parameters for a specific model

We did calculations for an extension of the *Standard Model* with strongly coupled scalar bosons to the Higgs field.

L. L., A. Megevand and A. D. Sanchez, JCAP **1210**, 024 (2012).

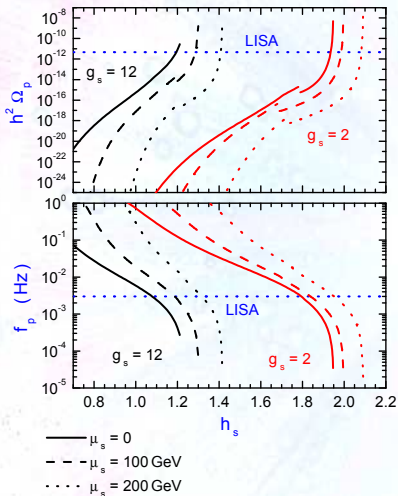


h_s is a coupling constant between scalars and Higgs
(It represents the strength of the interaction between these particles)

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Gravitational waves

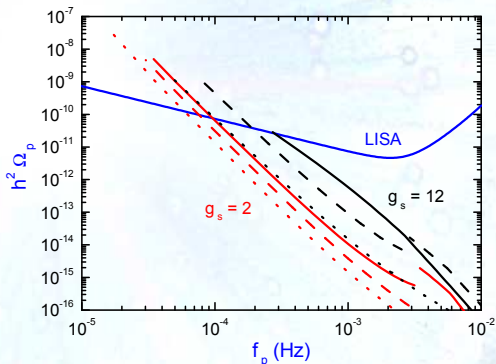
Extra coupled scalar bosons to the Higgs field.



$$h^2 \Omega_{\text{GW}}(f) \equiv \frac{h^2}{\rho_c} \frac{d\rho_{\text{GW}}}{d \log f}$$

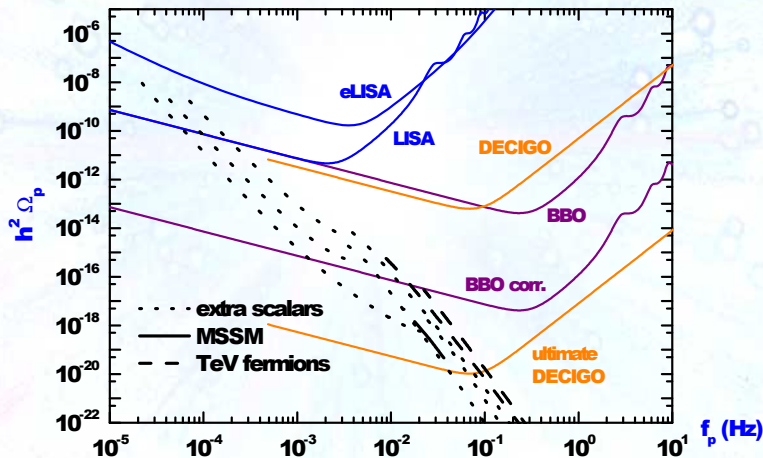
f_p = peak frequency


$h^2 \Omega_p = h^2 \Omega_{\text{GW}}(f_p)$ = peak intensity



Other models and detectors

It is difficult to get a signal observable by LISA, although we have seen that the conjunction between other models and planned detectors suggests greater chances of detection.



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The goal is to take into account the different bubble sizes present at bubble collisions (the literature assumes that the turbulence receives energy from a single length scale)

- Compute size distribution **OK**
- Compute turbulence (**Non-trivial**)
- Compute gravitational waves

Prospects

We have in mind:

- refine hydrodynamics calculations.
- consider alternative methods to detect the gravitational waves generated in cosmological phase transitions (CMB Polarization???)